

## A TRANSITION TO NON-RADIATING DIELECTRIC WAVEGUIDE.

J A G MALHERBE\*, J H CLOETE† AND I E LÖSCH\*

\*Department of Electronic Engineering and Laboratory for Advanced Engineering, University of Pretoria, Pretoria, 0002 South Africa.

†Previously at the above address; now at Department of Electrical and Electronic Engineering, University of Stellenbosch, Stellenbosch, 7600 South Africa.

## Abstract

A transition between standard X-band waveguide and Non-radiating Dielectric (NRD) waveguide is described. It features an impedance taper from air-to-dielectric-filled waveguide cascaded with an impedance taper and mode converter to NRD. In both cases the taper shape and length are obtained through a theoretical calculation. The main property of the transition is a low reflection coefficient so that conventional waveguide equipment can be used for measurements in NRD waveguide.

## INTRODUCTION

The Non-Radiating Dielectric Waveguide (NRD) has become increasingly more popular as more applications are being developed [1]-[4]. In order to enable the use of conventional waveguide equipment for measurements in NRD waveguide, high-quality transitions between NRD and standard waveguide are essential. This paper describes the development of such a transition, to match standard X-band waveguide (22.86x10.16mm) to NRD with a plate separation of 15.00 mm and dielectric with  $\epsilon_r = 2.55$  (Polypenco Q200.5 cross-linked polystyrene) of dimensions 15.00x10.16 mm. The design is based on an analytical procedure rather than an empirical method as described in [5].

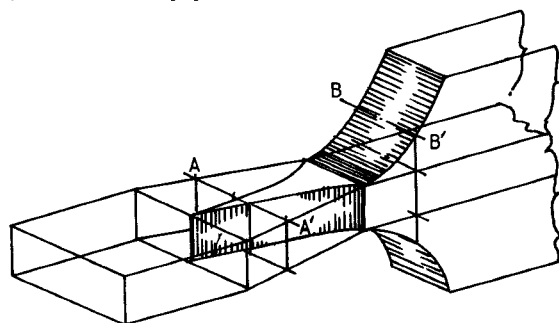


Fig 1. Composite Transition

The transition comprises two sections as shown in Fig 1; one is an impedance taper from air-to-dielectric-filled metal-walled waveguide. At all cross sections along the taper, this guide is as shown in Fig 2(a). The second section transforms from the reduced-width waveguide of 15.00x10.16 mm to NRD with 15.00 mm plate separation. Its cross section is shown in Fig 2(b) where the height,  $c$ , is a variable, and width,  $a$ , is constant.

In the section that follows, the properties of both

types of non-homogeneous guide will be briefly described. For each, it will be indicated that a taper can be designed that matches the guide properties at the extremities of each section, while performing the necessary mode conversion. Measured results will be described in conclusion.

## TRANSITION FROM AIR-TO DIELECTRIC-FILLED WAVEGUIDE

A typical cross section of this transition is shown in Fig 2(a). It consists of a linearly tapered metal waveguide of width  $c$ , partially filled with a dielectric of varying width  $a$ . It can be shown [6]

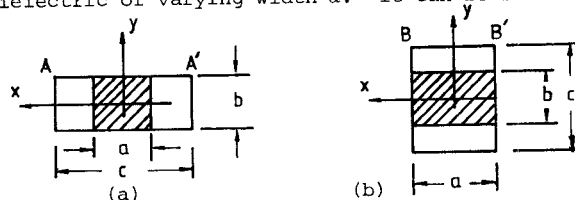


Fig 2. Non-homogeneous Cross Section that the fields in the dielectric-filled region are given by (with  $\exp(-j\beta_z z)$  suppressed in all cases),

$$E_{x_1} \equiv 0 \quad (1)$$

$$E_{y_1} = -j\omega \cos(\beta_{x_1} x) \quad (2)$$

$$E_{z_1} \equiv 0 \quad (3)$$

$$H_{x_1} = -\frac{j\beta_y}{\mu} \cos(\beta_{x_1} x) \quad (4)$$

$$H_{y_1} \equiv 0 \quad (5)$$

$$H_{z_1} = -\frac{\beta_{x_1}}{\mu} \sin(\beta_{x_1} x) \quad (6)$$

For the air-filled region,

$$E_{x_2} \equiv 0 \quad (7)$$

$$E_{y_2} = -j\omega K_1 \sin(\beta_{x_2} [a/2-x]) \quad (8)$$

$$E_{z_2} \equiv 0 \quad (9)$$

$$H_{x_2} = -j\omega K_1 \frac{\beta_g}{\mu} \sin(\beta_{x_2} [a/2-x]) \quad (10)$$

$$H_{y_2} \equiv 0 \quad (11)$$

$$H_{z_2} = -K_1 \frac{\beta_{x_2}}{\mu} \cos(\beta_{x_2} [a/2-x]) \quad (12)$$

where

$$K_1 = \cos(\beta_{x_1} c/2) \csc(\beta_{x_2} [a/2-c/2]) \quad (13)$$

and

$$\beta_{x_1} \tan(\beta_{x_1} c/2) = \beta_{x_2} \cot(\beta_{x_2} [a/2 - c/2]) \quad (14)$$

The transition tapers between values  $c_1 = 22.86$  mm,  $b = 10.16$  mm and  $a_1 = 0$ , and  $c_2 = 15.0$  mm,  $b = 10.16$  mm,  $a_2 = 15.0$  mm. For ease of manufacturing,

$$c(z) = c_1 + (c_2 - c_1)z/\ell \quad (15)$$

The equations (1) to (12) above indicate that a  $TE_{10}$  mode exists in the entire region. The dielectric must therefore have sides parallel to the metal walls. Hecken's impedance function [7] was selected for the calculation of the varying dielectric width, while the wave impedance was chosen as the appropriate impedance for calculations [8]. The material used as dielectric was Polypenco Q200.5 cross-linked polystyrene with  $\epsilon_r = 2.55$  and the start ( $z = 0$ ) and end ( $z = \ell$ ) parameters of the transition are shown in Table I.

TABLE I. WAVEGUIDE PARAMETERS AT EXTREMITIES OF TAPER

| Position                     | $z = 0$        | $z = \ell$     |
|------------------------------|----------------|----------------|
| Wave Impedance $Z$           | 550.4 $\Omega$ | 372.0 $\Omega$ |
| Guide wavelength $\lambda_g$ | 48.70 mm       | 32.92 mm       |
| Cutoff frequency $f_c$       | 6.56 GHz       | 6.65 GHz       |

The wave impedance and guide wavelength is computed from

$$Z = 376.7 \lambda_g / \lambda \quad (16)$$

$$\lambda_g = \lambda [1 - (\beta_{x_1}/k_0)^2]^{-1/2} \quad (17)$$

where  $k_0$  is the free space wave number, and  $\beta_{x_1}$  is found by numerically solving simultaneously equation (14) and

$$\beta_{x_2}^2 = \beta_{x_1}^2 - (\epsilon_r - 1)k_0^2 \quad (18)$$

The impedance function described by Hecken is for a TEM structure, where the variation of impedance,  $Z(\zeta)$  is calculated for *equispaced* points  $\zeta$ ; in this instance, because the guide wavelength varies as  $Z(z)$  varies, the function  $Z(\zeta)$  for *equispaced* points  $\zeta$  is converted to points *electrically*. The total taper length was chosen as  $\ell = 40$  mm; the impedance values calculated are shown in Table II.

TABLE II IMPEDANCES FOR -25dB HECKEN TAPER  
 $B = 3.2136 (\beta \ell_{\min}) = 4.1046$

| $\zeta$    | -1.0  | -0.8  | -0.6  | -0.4  | -0.2  | 0.0   |
|------------|-------|-------|-------|-------|-------|-------|
| $Z(\zeta)$ | 550.4 | 541.5 | 526.3 | 504.9 | 497.3 | 452.3 |
| $\zeta$    | 0.2   | 0.4   | 0.6   | 0.8   | 1.0   |       |
| $Z(\zeta)$ | 426.9 | 405.2 | 388.8 | 377.8 | 372.0 |       |

An iterative procedure was used with the impedance points *equispaced* as a starting value, and the guide wavelengths calculated. The spacing was adjusted and the process repeated until the points were *electrically* *equidistant*. Fig 3(a) shows the resulting taper to scale.

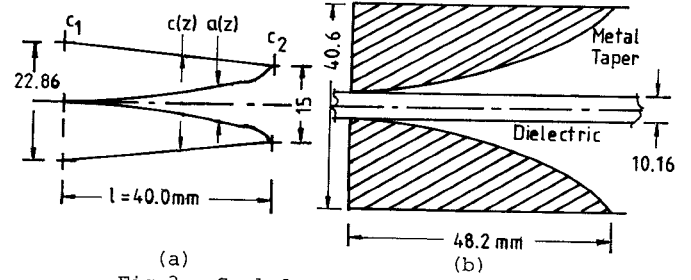


Fig 3. Scaled taper shapes  
TRANSITION FROM DIELECTRIC-FILLED WAVEGUIDE TO NRD

The cross section of this part of the transition is shown in Fig 2(b). The dielectric dimensions are kept constant at  $a = 15.0$  mm and  $b = 10.15$  mm, while the height is varied from  $c_1 = 15.0$  mm at  $z = 0$ , to infinity. In practice the NRD waveguide was made 40 mm high. The fields in this transitional region are given by [6], inside the dielectric,

$$E_{x_1} = \frac{(\pi/a)\beta_y}{j\omega\mu_0\epsilon_0\epsilon_r} \sin(\frac{\pi x}{a}) \sin(\beta_y y) \quad (19)$$

$$E_{y_1} = \frac{(k^2 - \beta_y^2)}{j\omega\mu_0\epsilon_0} \cos(\frac{\pi x}{a}) \cos(\beta_y y) \quad (20)$$

$$E_{z_1} = \frac{\beta_g \beta_y}{\omega\mu_0\epsilon_0\epsilon_r} \cos(\frac{\pi x}{a}) \sin(\beta_y y) \quad (21)$$

$$H_{x_1} = \frac{j\beta_g}{\mu_0} \cos(\frac{\pi x}{a}) \cos(\beta_y y) \quad (22)$$

$$H_{y_1} = 0 \quad (23)$$

$$H_{z_1} = -\frac{(\pi/a)}{\mu_0} \sin(\frac{\pi x}{a}) \cos(\beta_y y) \quad (24)$$

Note that  $\exp(-j\beta_z z)$  has again been suppressed. In the air-filled region,

$$E_{x_2} = K_2 \frac{(\pi/a)v}{j\omega\mu_0\epsilon_0} \sin(\pi x/a) \sinh[(c/2-y)v_y] \quad (25)$$

$$E_{y_2} = K_2 \frac{(k^2 + v^2)}{j\omega\mu_0\epsilon_0} \cos(\pi x/a) \cosh[(c/2-y)v_y] \quad (26)$$

$$E_{z_2} = K_2 \frac{\beta_g v}{\omega\mu_0\epsilon_0} \cos(\pi x/a) \sinh[(c/2-y)v_y] \quad (27)$$

$$H_{x_2} = jK_2 \frac{\beta_g}{\mu_0} \cos(\pi x/a) \cosh[(c/2-y)v_y] \quad (28)$$

$$H_{y_2} = 0 \quad (29)$$

$$H_{z_2} = -K_2 \frac{(\pi/a)}{\mu_0} \sin(\pi x/a) \cosh[(c/2-y)v_y] \quad (30)$$

where

$$K_2 = \cos(\beta_y \frac{b}{2}) \operatorname{sech}[(c/2-b/2)v_y] \quad (31)$$

and 
$$\tan(\beta \frac{b}{2}) \coth[a/2 - b/2] \frac{v_y}{y} = \frac{\epsilon_r v_y}{y} \quad (32)$$

The inherent mismatch of this section is 11.3dB, so that a Hecken impedance taper was designed to give an improvement of 30 dB. A similar procedure to that described in the previous paragraph was followed, using  $(\beta l_{\min}) = 4.7533$ ; 20 equal parts of electrical length were used and the scaled taper is shown in Fig 3(b).

Note that the limiting form of equations (19)-(32) above is, for  $c = b$ , a  $TE_{10}$ -mode in a dielectrically loaded waveguide. For  $c \rightarrow \infty$ , the equations describe the  $TM_0$ -mode of the NRD. Thus mode matching is simultaneously achieved.

#### MEASUREMENTS

A back-to-back set of the first type of taper was manufactured, with 250 mm of reduced width guide separating the two transitions. Fig 4 shows the measured reflection coefficient with the output

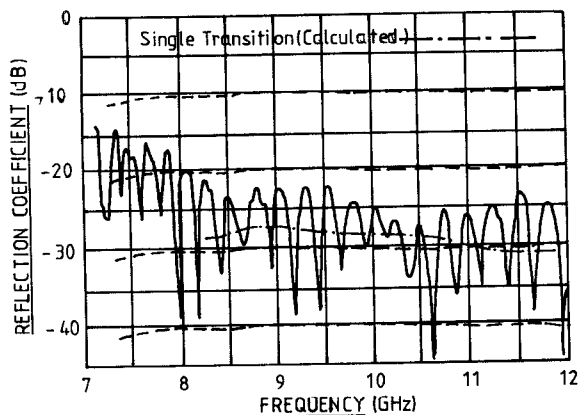


Fig 4. Measured performance of two back to back transitions between air- and dielectric-filled waveguide, separated by 250mm of dielectric filled guide.

transition terminated in a matched load. The oscillatory nature of the reflection coefficient is caused by the electrically varying length of line between the two transitions. The dotted line shows the calculated return loss of one transition; it is seen to be more than 27 dB over the useful frequency band of 8.5 to 10.5 GHz.

A complete transition was also manufactured. The NRD guide was terminated by laying two strips of Eccoshield FGM 40, each approximately 20 cm long and 15 mm wide on either side of the dielectric, with the leading edges of the absorber pulled away in the same shape as the launching taper. The measured reflection coefficient is never worse than -19 dB above 9.5 GHz, and is shown in Fig 5.

#### CONCLUSION

The design of a transition between standard waveguide and NRD has been described. The method is based on an analytical approach and the constructed model has shown excellent performance.

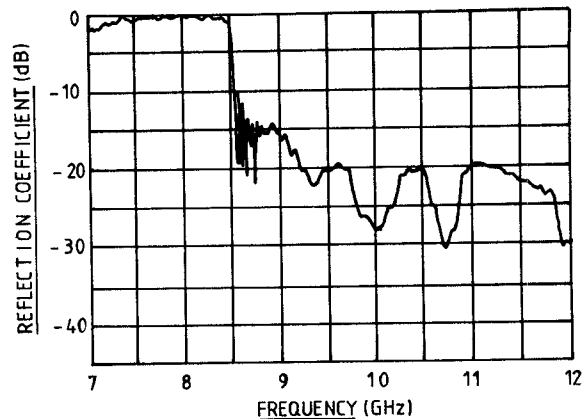


Fig 5. Measured performance of the composite transition between standard waveguide and terminated NRD.

#### ACKNOWLEDGMENT

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